



TITLE:

Construction of the least informative observable conserved by a given quantum instrument

AUTHOR(S):

Kuramochi, Yui

CITATION:

Kuramochi, Yui. Construction of the least informative observable conserved by a given quantum instrument. Journal of Mathematical Physics 2015, 56(9): 092202.

ISSUE DATE:

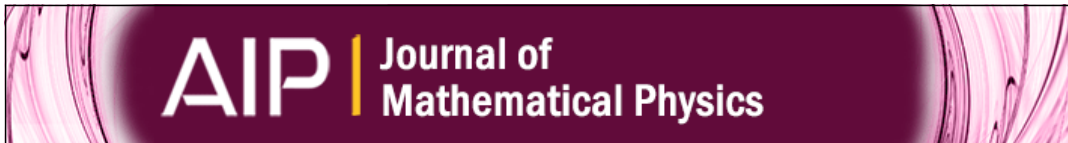
2015-09-28

URL:

<http://hdl.handle.net/2433/203028>

RIGHT:

© 2015 AIP Publishing LLC. This article may be downloaded for personal use only. Any other use requires prior permission of the author and the American Institute of Physics.



Construction of the least informative observable conserved by a given quantum instrument

Yui Kuramochi

Citation: *Journal of Mathematical Physics* **56**, 092202 (2015); doi: 10.1063/1.4931625

View online: <http://dx.doi.org/10.1063/1.4931625>

View Table of Contents: <http://scitation.aip.org/content/aip/journal/jmp/56/9?ver=pdfcov>

Published by the [AIP Publishing](#)

Articles you may be interested in

[Quantum Information](#)

Am. J. Phys. **78**, 654 (2010); 10.1119/1.3302835

[Energy conservation in quantum mechanics](#)

Am. J. Phys. **72**, 580 (2004); 10.1119/1.1648326

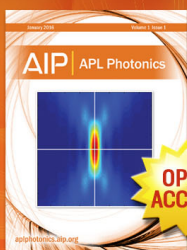
[Quantum Computation and Quantum Information](#)

Am. J. Phys. **70**, 558 (2002); 10.1119/1.1463744

[Conservation of information: Quantum mechanics and RIS](#)

AIP Conf. Proc. **584**, 87 (2001); 10.1063/1.1405588

[The qualitative theory of elementary transformations of one- and multichannel quantum systems in the inverse-problem approach. The construction of transformations with given spectral parameters](#)
Phys. Part. Nucl. **30**, 111 (1999); 10.1134/1.953100



Launching in 2016!

The future of applied photonics research is here

OPEN
ACCESS

AIP | APL
Photonics

Construction of the least informative observable conserved by a given quantum instrument

Yui Kuramochi^{a)}*Department of Nuclear Engineering, Kyoto University, 6158540 Kyoto, Japan*

(Received 27 May 2015; accepted 11 September 2015; published online 28 September 2015)

For a quantum measurement process described by a quantum instrument \mathcal{I} and a system observable corresponding to a positive-operator valued measure (POVM) E , \mathcal{I} is said to conserve the information of E if the joint successive measurement of \mathcal{I} followed by E is equivalent to a single measurement of E . We show that for any quantum instrument \mathcal{I} we can construct a POVM conserved by \mathcal{I} . Intuitively, the construction gives the infinite joint successive measurement of \mathcal{I} . We also show that the constructed POVM is the least informative observable among POVMs conserved by \mathcal{I} , i.e., the constructed POVM can be realized by a classical post-processing of any POVM conserved by \mathcal{I} . As typical examples of quantum instruments, we explicitly evaluate POVMs of infinite successive measurements for photon counting and quantum counter instruments. © 2015 AIP Publishing LLC.
[\[http://dx.doi.org/10.1063/1.4931625\]](http://dx.doi.org/10.1063/1.4931625)

I. INTRODUCTION

While quantum measurement back-action causes the state change and the information loss of the system, some kind of information is known to be conserved if we focus on a proper system observable described by a positive-operator valued measure (POVM). An example of such measurement processes is the photon counting measurement, also known as the Srinivas-Davies model,¹ in which we can estimate photon number of the pre-measurement state from that of the post-measurement state and the measurement outcome. In general, a measurement process described by a quantum instrument² \mathcal{I} is said to conserve a system observable described by a positive-operator valued measure E if the joint measurement of \mathcal{I} followed by E is equivalent to a single measurement of E .

In the previous work,³ the author examined some physical examples of quantum instruments and showed that some intuitively “natural” observable for each instrument satisfies the conservation condition. Then, it is natural to ask whether there exists a POVM E conserved by a given quantum instrument \mathcal{I} . The answer, the main result of this paper, is affirmative, and E can be constructed as the infinite successive joint measurement of \mathcal{I} , called an infinite composition of \mathcal{I} . Furthermore, this infinite successive measurement is shown to be characterized by the minimality up to the fuzzy preorder relation^{4–7} among the conserved POVMs, i.e., the infinite successive measurement is the least informative POVM conserved by \mathcal{I} . We also reconsider the photon counting and quantum counter instruments and explicitly derive their infinite compositions. As a by-product of the discussion, we also correct a mathematical insufficiency in the proof of the existing work⁸ concerning the convergence of the normalized count number in the quantum counter measurement.

In Ref. 3, the conservation condition was introduced as a sufficient condition for the “relative-entropy conservation law” for system observable E , which is the generalization of “Shannon entropy conservation” derived by Ban.⁹ In the comparison of two POVMs, the discussion of Ref. 3 is largely based on the concept of the sufficient statistic,¹⁰ while it can be formulated more rigorously and briefly by using the fuzzy preorder and equivalence relations between POVMs^{4–7} (Definition

^{a)}Electronic mail: kuramochi.yui.22c@st.kyoto-u.ac.jp

1). The reformulation of the information conservation condition based on the fuzzy equivalence relation is another purpose of the present paper.

This paper is organized as follows. In Sec. II, some preliminary results concerning the fuzzy preorder and equivalence relations and the composition of quantum measurement processes are reviewed. In Sec. III, we construct the infinite composition of a given quantum instrument with a standard Borel outcome space and show that it is the least informative POVM that is conserved by the instrument. In Sec. IV, we consider photon counting and quantum counter instruments and derive the explicit forms of their infinite compositions. Sec. V summarizes the main results of this paper.

II. PRELIMINARIES

In this section we briefly review some preliminary results on the quantum theory of measurement and fix the notation. For a general reference of quantum measurement, we refer Refs. 5, 11, and 12.

A. Positive-operator valued measures and fuzzy preorder and equivalence relations

We fix a complex Hilbert space \mathcal{H} and denote the set of bounded linear operators on \mathcal{H} as $\mathcal{L}(\mathcal{H})$. Let (Ω, \mathcal{B}) be a measurable space. A mapping $E : \mathcal{B} \rightarrow \mathcal{L}(\mathcal{H})$ is called a POVM with its outcome space (Ω, \mathcal{B}) if

- (i) $E(B) \geq 0$ ($\forall B \in \mathcal{B}$),
- (ii) $E(\Omega) = I$,
- (iii) for any disjoint $\{B_i\}_{i \geq 1} \subset \mathcal{B}$, $E(\cup_{i \geq 1} B_i) = \sum_{i \geq 1} E(B_i)$ in the weak operator topology.

Here O and I are zero and identity operators, respectively. Let E be a POVM with its outcome space (Ω, \mathcal{B}) . A measurable set $N \in \mathcal{B}$ is called an E -null set if $E(N) = O$. E -almost sure equations and convergences for stochastic variables are also defined in a similar manner as for a classical probability measure.

Let $(\Omega_1, \mathcal{B}_1)$ and $(\Omega_2, \mathcal{B}_2)$ be measurable spaces. A mapping $\nu(\cdot) : \Omega_2 \times \mathcal{B}_1 \ni (\omega_2, B) \mapsto \nu_{\omega_2}(B) \in [0, 1]$ is said to be a *Markov kernel* if

- (i) for each $\omega_2 \in \Omega_2$, the mapping $\nu_{\omega_2}(\cdot) : \mathcal{B}_1 \rightarrow [0, 1]$ is a probability measure,
- (ii) for each $B \in \mathcal{B}_1$, the mapping $\nu(B) : \Omega_2 \rightarrow [0, 1]$ is \mathcal{B}_2 -measurable.

Let E^1 and E^2 be POVMs with their outcome spaces $(\Omega_1, \mathcal{B}_1)$ and $(\Omega_2, \mathcal{B}_2)$, respectively. Following Refs. 5–7, we define a relation $E^1 \leq E^2$ by the existence of a Markov kernel $\nu(\cdot) : \Omega_2 \times \mathcal{B}_1 \rightarrow [0, 1]$ such that

$$E^1(B) = \int_{\Omega_2} \nu_{\omega_2}(B) E^2(d\omega_2) \quad (\forall B \in \mathcal{B}_1), \quad (1)$$

and we say that E^1 is *fuzzier* than E^2 . POVMs E^1 and E^2 are said to be *equivalent*, denoted as $E^1 \simeq E^2$, if $E^1 \leq E^2$ and $E^2 \leq E^1$. The relations \leq and \simeq are preorder and equivalence relations for POVMs,^{6,7} respectively. Equation (1) intuitively means that the measurement of E^1 can be realized by performing the measurement of E^2 and the classical information processing on the measurement outcome ω_2 .

The following lemma concerning the Markov kernel will be used later.

Lemma 1. 1. Let $(\Omega_i, \mathcal{B}_i)$ ($i = 1, 2, 3$) be measurable spaces. Suppose that

$$\begin{aligned} \nu^1(\cdot) : \Omega_2 \times \mathcal{B}_1 &\rightarrow [0, 1], \\ \nu^2(\cdot) : \Omega_3 \times \mathcal{B}_2 &\rightarrow [0, 1] \end{aligned}$$

are Markov kernels. Then $\nu^3(\cdot)$ defined by

$$\nu_{\omega_3}^3(B) := \int_{\Omega_2} \nu_{\omega_2}^1(B) \nu_{\omega_3}^2(d\omega_2) \quad (B \in \mathcal{B}_1, \omega_3 \in \Omega_3)$$

is a Markov kernel.

2. Let $(\Omega_i, \mathcal{B}_i)$ ($i = 1, 2, 3$) be measurable spaces and let $\nu(\cdot): \Omega_3 \times \mathcal{B}_2 \rightarrow [0, 1]$ be a Markov kernel. Define $\tilde{\nu}(\cdot)$ by

$$\begin{aligned}\tilde{\nu}_{(\omega_1, \omega_3)}(B) &:= \nu_{\omega_3}(B|_{\omega_1}), \\ B|_{\omega_1} &:= \{\omega_2 \in \Omega_2 | (\omega_1, \omega_2) \in B\}\end{aligned}\quad (2)$$

for $(\omega_1, \omega_3) \in \Omega_1 \times \Omega_3$ and $B \in \mathcal{B}_1 \times \mathcal{B}_2$. Then $\tilde{\nu}(\cdot)$ is a Markov kernel.

Proof. 1 is shown in the proof of Proposition 1 in Ref. 6. We show 2. For each $(\omega_1, \omega_3) \in \Omega_1 \times \Omega_3$, $\tilde{\nu}_{(\omega_1, \omega_3)}(\cdot)$ is a probability measure. To show the measurability of $\tilde{\nu}(B)$ ($B \in \mathcal{B}_1 \times \mathcal{B}_2$), define a class \mathcal{D} of subsets of $\Omega_1 \times \Omega_2$ by

$$\mathcal{D} := \{B \in \mathcal{B}_1 \times \mathcal{B}_2 | \tilde{\nu}(B) \text{ is } \mathcal{B}_1 \times \mathcal{B}_3\text{-measurable}\}.$$

Then \mathcal{D} is a Dynkin system (λ -system), i.e., it is closed under countable disjoint unions and proper differences and contains $\Omega_1 \times \Omega_2$. For each $B_i \in \mathcal{B}_i$ ($i = 1, 2$),

$$\tilde{\nu}_{(\omega_1, \omega_3)}(B_1 \times B_2) = \chi_{B_1}(\omega_1) \nu_{\omega_3}(B_2)$$

is $\mathcal{B}_1 \times \mathcal{B}_3$ -measurable, where $\chi_B(\cdot)$ is an indicator function for a subset B . Thus, \mathcal{D} contains the class of the cylinder sets $\{B_1 \times B_2 | B_1 \in \mathcal{B}_1, B_2 \in \mathcal{B}_2\}$ and the Dynkin's theorem assures that \mathcal{D} coincides with $\mathcal{B}_1 \times \mathcal{B}_2$, which proves the assertion. \square

B. Completely positive instruments and their compositions

Let (Ω, \mathcal{B}) be a measurable space. A *completely positive (CP) instrument*^{2,12,13} (in the Heisenberg picture) with its outcome space (Ω, \mathcal{B}) is a mapping

$$\mathcal{I}(\cdot): \mathcal{B} \times \mathcal{L}(\mathcal{H}) \ni (B, a) \mapsto \mathcal{I}_B(a) \in \mathcal{L}(\mathcal{H})$$

such that

- (i) for any $B \in \mathcal{B}$, $\mathcal{I}_B(\cdot): \mathcal{L}(\mathcal{H}) \rightarrow \mathcal{L}(\mathcal{H})$ is a normal CP linear map,
- (ii) for any disjoint $\{B_i\}_{i \geq 1} \subset \mathcal{B}$ and any $a \in \mathcal{L}(\mathcal{H})$, $\mathcal{I}_{\cup_{i \geq 1} B_i}(a) = \sum_{i \geq 1} \mathcal{I}_{B_i}(a)$ in the ultraweak operator topology,
- (iii) $\mathcal{I}_\Omega(I) = I$.

We also define a positive (P) instrument by replacing CP with P in the definition of the CP instrument. A CP instrument $\mathcal{I}(\cdot)$ with its outcome space (Ω, \mathcal{B}) describes the statistics of the measurement outcome and the state change due to the measurement simultaneously in a necessary and sufficient manner. The POVM corresponding to the probability distribution of the measurement outcome is given by $E(B) = \mathcal{I}_B(I)$ ($B \in \mathcal{B}$). Here we only consider the case when the range and the domain of $\mathcal{I}_B(\cdot)$ are identical, which is a necessary condition so that we can compare the *same* observable of the system before and after the measurement.

A measurable space (Ω, \mathcal{B}) Borel isomorphic to a complete separable metric space is called a *standard Borel space*.¹⁴ A (C)P instrument (respectively, POVM) with a standard Borel outcome space is called a standard Borel (C)P instrument (respectively, a standard Borel POVM). *In the rest of this paper, we only consider standard Borel (C)P instruments and POVMs.*

The following theorem, which is a slight modification of the theorem due to Davies and Lewis,^{12,13} assures the existence of a CP instrument and a POVM corresponding to a joint successive measurement process.

Theorem 1. (i) Let $\mathcal{I}^i(\cdot)$ be a CP instrument with a standard Borel outcome space $(\Omega_i, \mathcal{B}_i)$ ($i = 1, 2$). Then there exists a unique CP instrument $\mathcal{I}^{12}(\cdot)$ with the product outcome space $(\Omega_1 \times \Omega_2, \mathcal{B}_1 \times \mathcal{B}_2)$ such that

$$\mathcal{I}_{B_1 \times B_2}^{12}(\cdot) = \mathcal{I}_{B_1}^1 \circ \mathcal{I}_{B_2}^2(\cdot) \quad (B_1 \in \mathcal{B}_1, B_2 \in \mathcal{B}_2). \quad (3)$$

- (ii) Let $\mathcal{I}^1(\cdot)$ be a CP instrument with a standard Borel outcome space $(\Omega_1, \mathcal{B}_1)$ and let E^2 be a POVM with a standard Borel outcome space $(\Omega_2, \mathcal{B}_2)$. Then there exists a unique POVM E^{12}

with the product outcome space $(\Omega_1 \times \Omega_2, \mathcal{B}_1 \times \mathcal{B}_2)$ such that

$$E^{12}(B_1 \times B_2) = \mathcal{I}_{B_1}^1(E^2(B_2)) \quad (B_1 \in \mathcal{B}_1, B_2 \in \mathcal{B}_2). \quad (4)$$

We call the CP instrument $\mathcal{I}^{12}(\cdot)$ and the POVM E^{12} as *compositions*, and denote them as $(\mathcal{I}^1 * \mathcal{I}^2)(\cdot)$ and $\mathcal{I}^1 * E^2$, respectively.

Proof. We first show (i). According to Theorem 4.2.2 of Ref. 12, there exists a unique P instrument $\mathcal{I}^{12}(\cdot)$ such that condition (3) holds. To show the complete positivity of $\mathcal{I}^{12}(\cdot)$, define a class \mathcal{D} of subsets of $\Omega_1 \times \Omega_2$ by

$$\mathcal{D} := \{B \in \mathcal{B}_1 \times \mathcal{B}_2 \mid \mathcal{I}_B^{12}(\cdot) \text{ is CP}\}.$$

From condition (3), \mathcal{D} contains the class of cylinder sets. Since we can easily verify that \mathcal{D} is a Dynkin class, Dynkin's theorem assures that $\mathcal{D} = \mathcal{B}_1 \times \mathcal{B}_2$, proving assertion (i).

To show (ii), take a CP instrument $\mathcal{I}^2(\cdot)$ with the outcome space $(\Omega_2, \mathcal{B}_2)$ such that $\mathcal{I}_B^2(I) = E^2(B)$ ($B \in \mathcal{B}_2$). Then a POVM $E^{12}(\cdot) := (\mathcal{I}^1 * \mathcal{I}^2)(I)$ satisfies condition (4). The uniqueness can be shown by using the Dynkin's theorem as parallel as in the classical measure. \square

Now let $\mathcal{I}^k(\cdot)$ be a CP instrument with a standard Borel outcome space $(\Omega_k, \mathcal{B}_k)$ ($k = 1, 2, \dots$) and let E be a POVM with a standard Borel outcome space (Ω, \mathcal{B}) . Then we have the following associative laws for the composition:

$$\begin{aligned} (\mathcal{I}^1 * \mathcal{I}^2) * \mathcal{I}^3 &= \mathcal{I}^1 * (\mathcal{I}^2 * \mathcal{I}^3), \\ (\mathcal{I}^1 * \mathcal{I}^2) * E &= \mathcal{I}^1 * (\mathcal{I}^2 * E). \end{aligned}$$

Thus, we may write these CP instrument and POVM as $\mathcal{I}^1 * \mathcal{I}^2 * \mathcal{I}^3$ and $\mathcal{I}^1 * \mathcal{I}^2 * E$, respectively. Multiple compositions $\mathcal{I}^1 * \dots * \mathcal{I}^n$ and $\mathcal{I}^1 * \dots * \mathcal{I}^n * E$ for general $n \geq 1$ are also defined in a similar manner. These are the unique CP instrument and POVM such that

$$\begin{aligned} (\mathcal{I}^1 * \dots * \mathcal{I}^n)_{B_1 \times \dots \times B_n}(\cdot) &= \mathcal{I}_{B_1}^1 \circ \dots \circ \mathcal{I}_{B_n}^n(\cdot), \\ (\mathcal{I}^1 * \dots * \mathcal{I}^n * E)(B_1 \times \dots \times B_n \times B) &= \mathcal{I}_{B_1}^1 \circ \dots \circ \mathcal{I}_{B_n}^n(E(B)) \end{aligned}$$

for each $B_k \in \mathcal{B}_k$ ($k = 1, \dots, n$) and $B \in \mathcal{B}$.

For later use, we show the following lemmas.

Lemma 2. Let $\mathcal{I}^1(\cdot)$ be a CP instrument with a standard Borel outcome space $(\Omega_1, \mathcal{B}_1)$ and let E^2 be a POVM with a standard Borel outcome space $(\Omega_2, \mathcal{B}_2)$. Then for each $B_1 \in \mathcal{B}_1$ and for each bounded complex valued \mathcal{B}_2 -measurable function f ,

$$\int_{\Omega_1 \times \Omega_2} \chi_{B_1}(\omega_1) f(\omega_2) (\mathcal{I}^1 * E^2)(d\omega_1 \times d\omega_2) = \mathcal{I}_{B_1}^1 \left(\int_{\Omega_2} f(\omega_2) E^2(d\omega_2) \right). \quad (5)$$

Proof. It is sufficient to show Eq. (5) when $f \geq 0$. When f is a measurable simple function, Eq. (5) holds. For general f , take a monotone sequence of non-negative measurable simple functions f_n such that $f_n(\omega_2) \uparrow f(\omega_2)$ for each $\omega_2 \in \Omega_2$. Then from the dominated convergence theorem we have

$$\int_{\Omega_1 \times \Omega_2} \chi_{B_1}(\omega_1) f_n(\omega_2) (\mathcal{I}^1 * E^2)(d\omega_1 \times d\omega_2) \uparrow \int_{\Omega_1 \times \Omega_2} \chi_{B_1}(\omega_1) f(\omega_2) (\mathcal{I}^1 * E^2)(d\omega_1 \times d\omega_2), \quad (6)$$

$$\int_{\Omega_2} f_n(\omega_2) E^2(d\omega_2) \uparrow \int_{\Omega_2} f(\omega_2) E^2(d\omega_2). \quad (7)$$

Since $\mathcal{I}_{B_1}^1(\cdot)$ is normal, Eq. (7) implies that

$$\mathcal{I}_{B_1}^1 \left(\int_{\Omega_2} f_n(\omega_2) E^2(d\omega_2) \right) \uparrow \mathcal{I}_{B_1}^1 \left(\int_{\Omega_2} f(\omega_2) E^2(d\omega_2) \right). \quad (8)$$

Since LHSs of Eqs. (6) and (8) coincide, we obtain Eq. (5). \square

Lemma 3. Let $\mathcal{I}^1(\cdot)$ be a CP instrument with a standard Borel outcome space $(\Omega_1, \mathcal{B}_1)$ and let E^2 and E^3 be POVMs with standard Borel outcome spaces $(\Omega_2, \mathcal{B}_2)$ and $(\Omega_3, \mathcal{B}_3)$, respectively. Then we have

- (i) if $E^2 \leq E^3$, then $\mathcal{I}^1 * E^2 \leq \mathcal{I}^1 * E^3$,
- (ii) if $E^2 \simeq E^3$, then $\mathcal{I}^1 * E^2 \simeq \mathcal{I}^1 * E^3$.

Proof. We first show (i). From the assumption $E^2 \leq E^3$, there exists a Markov kernel $\nu(\cdot): \Omega_3 \times \mathcal{B}_2 \rightarrow [0, 1]$ such that

$$E^2(B_2) = \int_{\Omega_3} \nu_{\omega_3}(B_2) E^3(d\omega_3) \quad (B_2 \in \mathcal{B}_2).$$

From Lemma 1, we can define a POVM E^{12} by

$$E^{12}(B) := \int_{\Omega_1 \times \Omega_3} \nu_{\omega_3}(B|_{\omega_1})(\mathcal{I}^1 * E^3)(d\omega_1 \times d\omega_3) \quad (B \in \mathcal{B}_1 \times \mathcal{B}_2).$$

From the definition of E^{12} , $E^{12} \leq \mathcal{I}^1 * E^3$ holds. On the other hand, for each $B_1 \in \mathcal{B}_1$ and $B_2 \in \mathcal{B}_2$, we have

$$\begin{aligned} E^{12}(B_1 \times B_2) &= \int_{\Omega_1 \times \Omega_3} \chi_{B_1}(\omega_1) \nu_{\omega_3}(B_2) (\mathcal{I}^1 * E^3)(d\omega_1 \times d\omega_3) \\ &= \mathcal{I}_{B_1}^1 \left(\int_{\Omega_3} \nu_{\omega_3}(B_2) E^3(d\omega_3) \right) \\ &= \mathcal{I}_{B_1}^1(E^2(B_2)) = (\mathcal{I}^1 * E^2)(B_1 \times B_2), \end{aligned}$$

where we have used Lemma 2 in the derivation of the second equality. Therefore, we obtain $\mathcal{I}^1 * E^2 = E^{12} \leq \mathcal{I}^1 * E^3$ and (i) is proved. (ii) immediately follows from (i) and the definition of the equivalence relation \simeq . \square

III. INFINITE COMPOSITION OF AN INSTRUMENT AND ITS MINIMAL INFORMATION-CONSERVING PROPERTY

Definition 1. Let $\mathcal{I}(\cdot)$ be a standard Borel CP instrument and let E be a standard Borel POVM. We say that E is conserved by \mathcal{I} , or \mathcal{I} conserves E , if $\mathcal{I} * E \simeq E$.

This condition is essentially the same as the one obtained in the author's previous work³ for a sufficient condition for the “relative-entropy conservation law” which is a generalization of the “Shannon-entropy conservation law” derived by Ban.⁹

From Definition 1 and Lemma 3, we immediately obtain the following theorem.

Theorem 2. Let $\mathcal{I}(\cdot)$ be a standard Borel CP instrument and let E^1 and E^2 be standard Borel POVMs. Suppose that $E^1 \simeq E^2$. Then $\mathcal{I}(\cdot)$ conserves E^1 if and only if $\mathcal{I}(\cdot)$ conserves E^2 . In other words, the conservation by $\mathcal{I}(\cdot)$ is well-defined to \simeq -equivalence classes of standard Borel POVMs.

Now, we ask whether there exists a standard Borel POVM E that is conserved by a given standard Borel CP instrument $\mathcal{I}(\cdot)$. The answer is affirmative, and it is given by a POVM called the *infinite composition* of \mathcal{I} corresponding to an infinite successive measurement of \mathcal{I} . The conservation of the infinite composition E_∞ by a CP instrument \mathcal{I} is intuitively understood as follows: the composition $\mathcal{I} * E_\infty$ is a measurement process in which we first perform \mathcal{I} and then perform \mathcal{I} infinitely many times, which is obviously equivalent to performing E_∞ . Furthermore, we will show in Theorem 4 that the infinite composition E_∞ of \mathcal{I} is special to \mathcal{I} in the sense that E_∞ is the minimal element with respect to the preorder relation \leq among POVMs conserved by \mathcal{I} .

The following proposition due to Tumulka¹⁵ is a key to the construction of the infinite composition, which is a POVM version of the celebrated Kolmogorov extension theorem for probability measures.¹⁶

Proposition 1 (quantum Kolmogorov extension theorem). Let $(\Omega_i, \mathcal{B}_i)$ ($i = 1, 2, \dots$) be a standard Borel space and let E_n ($n = 1, 2, \dots$) be a POVM with the product outcome space $(\prod_{i=1}^n \Omega_i, \prod_{i=1}^n \mathcal{B}_i)$. Suppose that $\{E_n\}$ satisfies the condition

$$E_n(B) = E_{n+1}(B \times \Omega_{n+1}) \quad \left(n \geq 1, B \in \prod_{i=1}^n \mathcal{B}_i \right). \quad (9)$$

Then, there exists a unique POVM E_∞ with the infinite product outcome space $(\prod_{i=1}^\infty \Omega_i, \prod_{i=1}^\infty \mathcal{B}_i)$ such that

$$E_n(B) = E_\infty \left(B \times \prod_{i=n+1}^\infty \Omega_i \right) \quad \left(n \geq 1, B \in \prod_{i=1}^n \mathcal{B}_i \right). \quad (10)$$

Condition (9) is called a *Kolmogorov consistency condition* and the POVM E_∞ satisfying (10) is said to be *consistent* with $\{E_n\}$. Note that the consistent POVM E_∞ is a standard Borel POVM since a countable product of standard Borel spaces is also a standard Borel space.

For simplicity, if $(\Omega_i, \mathcal{B}_i)$ ($1 \leq i \leq n$) is identical to (Ω, \mathcal{B}) , the product space $(\prod_{i=1}^n \Omega_i, \prod_{i=1}^n \mathcal{B}_i)$ ($1 \leq n \leq \infty$) is denoted as $(\Omega^n, \mathcal{B}^n)$. We also denote $\underbrace{\mathcal{I} * \dots * \mathcal{I}}_{n \text{ elements}}$ as \mathcal{I}^{*n} for a CP instrument \mathcal{I} .

Now we construct the infinite composition of a CP instrument.

Theorem 3 (infinite composition of an instrument). Let $\mathcal{I}(\cdot)$ be a CP instrument with a standard Borel outcome space (Ω, \mathcal{B}) . Then there exists a unique POVM E_∞ with the infinite product outcome space $(\Omega^\infty, \mathcal{B}^\infty)$ such that

$$E_\infty \left(\prod_{i=1}^n B_i \times \Omega^\infty \right) = \mathcal{I}_{B_1} \circ \dots \circ \mathcal{I}_{B_n}(\mathcal{I}) \quad (11)$$

for each $n \geq 1$ and $B_i \in \mathcal{B}$ ($i = 1, \dots, n$). The POVM E_∞ is called an *infinite composition* of \mathcal{I} .

Proof. Let us define a POVM $E_n(\cdot) := (\mathcal{I}^{*n})(\mathcal{I})$ for each $n \geq 1$. Since $E_n(\prod_{i=1}^n B_i)$ coincides with the RHS of Eq. (11), from the quantum Kolmogorov extension theorem, it is sufficient to show that $\{E_n\}$ satisfies the Kolmogorov consistency condition. For each $B \in \mathcal{B}^n$ ($n \geq 1$), we have

$$\begin{aligned} E_{n+1}(B \times \Omega) &= (\mathcal{I}^{*n} * \mathcal{I})_{B \times \Omega}(\mathcal{I}) \\ &= \mathcal{I}_B^{*n}(\mathcal{I}_\Omega(\mathcal{I})) \\ &= \mathcal{I}_B^{*n}(\mathcal{I}) = E_n(B), \end{aligned}$$

and the theorem holds. \square

The next theorem is the main result of this paper, which states that for a given CP instrument \mathcal{I} , the infinite composition of \mathcal{I} is the least informative POVM conserved by \mathcal{I} .

Theorem 4. Let $\mathcal{I}(\cdot)$ be a CP instrument with a standard Borel outcome space (Ω, \mathcal{B}) and let E_∞ be the infinite composition of \mathcal{I} . Then, \mathcal{I} conserves E_∞ . Furthermore, E_∞ is the minimal element with respect to the preorder relation \leq among the standard Borel POVMs conserved by \mathcal{I} .

Proof. $(\Omega^\infty, \mathcal{B}^\infty)$ and $(\Omega \times \Omega^\infty, \mathcal{B} \times \mathcal{B}^\infty)$ are Borel isomorphic by the mapping

$$\Omega \times \Omega^\infty \ni \left(\omega, \prod_{i=1}^\infty \omega_i \right) \mapsto (\omega, \omega_1, \omega_2, \dots) \in \Omega^\infty,$$

and we consider $\mathcal{I} * E_\infty$ as a POVM with the outcome space $(\Omega^\infty, \mathcal{B}^\infty)$ by this identification. Then for each $B, B_1, \dots, B_n \in \mathcal{B}$ ($1 \leq n < \infty$), we have

$$\begin{aligned} (\mathcal{I} * E_\infty)(B \times B_1 \times \dots \times B_n \times \Omega^\infty) &= \mathcal{I}_B(E_\infty(B_1 \times \dots \times B_n \times \Omega^\infty)) \\ &= \mathcal{I}_B \circ \mathcal{I}_{B_1} \circ \dots \circ \mathcal{I}_{B_n}(\mathcal{I}) \\ &= E_\infty(B \times B_1 \times \dots \times B_n \times \Omega^\infty). \end{aligned}$$

From the uniqueness of E_∞ , we obtain $\mathcal{I} * E_\infty = E_\infty$, which proves the conservation of E_∞ by \mathcal{I} .

To show the minimality of E_∞ , take an arbitrary POVM F with a standard Borel outcome space $(\Omega_X, \mathcal{B}_X)$ such that $F \simeq \mathcal{I} * F$. The goal of the proof is to construct a Markov kernel corresponding to an information-processing from Ω_X to Ω^∞ . Since $\mathcal{I} * F \leq F$, there exists a Markov kernel $\tilde{\nu}^1(\cdot): \Omega_X \times (\mathcal{B} \times \mathcal{B}_X) \rightarrow [0, 1]$ such that

$$(\mathcal{I} * F)(B) = \int_{\Omega_X} \tilde{\nu}_x^1(B) F(dx)$$

for each $B \in \mathcal{B} \times \mathcal{B}_X$. Since the Markov kernel $\tilde{\nu}_x^1(d\omega_1 \times dx_1)$ corresponds to a classical information processing generating a measurement outcome (ω_1, x_1) of $\mathcal{I} * F$ from a given measurement outcome x of F , we can construct a new Markov kernel $\tilde{\nu}_x^2(d\omega_1 \times d\omega_2 \times dx_2)$ which generates the measurement outcome of $\mathcal{I} * \mathcal{I} * F$ by applying the same classical information processing $\tilde{\nu}_{x_1}^1(d\omega_2 \times dx_2)$ to x_1 which generates (ω_2, x_2) . Repeating the same discussion, we can construct a sequence of Markov kernels $\tilde{\nu}_x^n(d\omega_1 \times \cdots \times d\omega_n \times dx_n)$ that generates the measurement outcome of $\mathcal{I}^{*n} * F$ from that of F . The formal definition of the sequence $\{\tilde{\nu}^n(\cdot)\}$ is given by

$$\tilde{\nu}_x^{n+1}(B) := \int_{\Omega^n \times \Omega_X} \tilde{\nu}_{x_n}^1(B|_{\omega^{(n)}}) \tilde{\nu}_x^n(d\omega^{(n)} \times dx_n) \quad (12)$$

for each $1 \leq n < \infty$ and each $B \in \mathcal{B}^{n+1} \times \mathcal{B}_X$, where we denote $\prod_{i=1}^n \omega_i$ as $\omega^{(n)}$ and $B|_{\omega^{(n)}} := \{(\omega_{n+1}, x_{n+1}) \in \Omega \times \Omega_X | (\omega^{(n)}, \omega_{n+1}, x_{n+1}) \in B\} (1 \leq n \leq \infty)$. From Lemma 1, $\tilde{\nu}^n(\cdot)$ defined by Eq. (12) is a well-defined Markov kernel. Now we show that

$$(\mathcal{I}^{*n} * F)(\cdot) = \int_{\Omega_X} \tilde{\nu}_x^n(\cdot) F(dx) \quad (13)$$

for each $n \geq 1$. If $n = 1$, Eq. (13) is evident from the definition of $\tilde{\nu}^1(\cdot)$. If Eq. (13) holds for $n \geq 1$, then for each $B_n \in \mathcal{B}^n, B_1 \in \mathcal{B}$, and $B_X \in \mathcal{B}_X$, we obtain

$$\begin{aligned} & \int_{\Omega_X} \tilde{\nu}_x^{n+1}(B_n \times B_1 \times B_X) F(dx) \\ &= \int_{\Omega_X} \left(\int_{\Omega^n \times \Omega_X} \tilde{\nu}_{x_n}^1((B_n \times B_1 \times B_X)|_{\omega^{(n)}}) \tilde{\nu}_x^n(d\omega^{(n)} \times dx_n) \right) F(dx) \\ &= \int_{\Omega_X} \left(\int_{\Omega^n \times \Omega_X} \chi_{B_n}(\omega^{(n)}) \tilde{\nu}_{x_n}^1(B_1 \times B_X) \tilde{\nu}_x^n(d\omega^{(n)} \times dx_n) \right) F(dx) \\ &= \int_{\Omega^n \times \Omega_X} \chi_{B_n}(\omega^{(n)}) \tilde{\nu}_{x_n}^1(B_1 \times B_X) (\mathcal{I}^{*n} * F)(d\omega^{(n)} \times dx_n) \end{aligned} \quad (14)$$

$$\begin{aligned} &= (\mathcal{I}^{*n})_{B_n} \left(\int_{\Omega_X} \tilde{\nu}_{x_n}^1(B_1 \times B_X) F(dx_n) \right) \\ &= (\mathcal{I}^{*(n+1)} * F)(B_n \times B_1 \times B_X), \end{aligned} \quad (15)$$

where in deriving Eqs. (15) and (14) we have used Lemma 2 and an equality

$$\begin{aligned} & \int_{\Omega_X} \left(\int_{\Omega^n \times \Omega_X} f(\omega^{(n)}, x_n) \tilde{\nu}_x^n(d\omega^{(n)} \times dx_n) \right) F(dx) \\ &= \int_{\Omega^n \times \Omega_X} f(\omega^{(n)}, x_n) (\mathcal{I}^{*n} * F)(d\omega^{(n)} \times dx_n) \end{aligned} \quad (16)$$

valid for any $\mathcal{B}^n \times \mathcal{B}_X$ -measurable bounded function f . Equation (16) holds when f is a simple function from the assumption of the induction, and for general f , we can prove the equation by taking a monotone sequence of simple functions converging pointwise to f . Thus we have shown Eq. (13)

Now we define a sequence of Markov kernels $\{\nu^n(\cdot)\}$ by

$$\nu_x^n(B) := \tilde{\nu}_x^n(B \times \Omega_X) \quad (x \in \Omega_X, B \in \mathcal{B}^n).$$

From the definition of $\tilde{\nu}^n(\cdot)$, for each $B \in \mathcal{B}^n$ we have

$$\begin{aligned}\nu_x^{n+1}(B \times \Omega) &= \int_{\Omega^n \times \Omega_X} \tilde{\nu}_{x_n}^1((B \times \Omega \times \Omega_X)|_{\omega(n)}) \tilde{\nu}_x^n(d\omega^{(n)} \times dx_n) \\ &= \int_{\Omega^n \times \Omega_X} \chi_B(\omega^{(n)}) \tilde{\nu}_x^n(d\omega^{(n)} \times dx_n) \\ &= \tilde{\nu}_x^n(B \times \Omega_X) = \nu_x^n(B).\end{aligned}$$

Thus, from Kolmogorov extension theorem, there exists a probability measure $\nu_x^\infty(\cdot)$ with the outcome space $(\Omega^\infty, \mathcal{B}^\infty)$ such that

$$\nu_x^\infty(B \times \Omega^\infty) = \nu_x^n(B) \quad (n \geq 1, B \in \mathcal{B}^n) \quad (17)$$

for each $x \in \Omega_X$. To show the \mathcal{B}_X -measurability of $\nu_x^\infty(B)$ ($B \in \mathcal{B}^\infty$), define a class \mathcal{D} of subsets of Ω^∞ by

$$\mathcal{D} := \{B \in \mathcal{B}^\infty | \nu_x^\infty(B) \text{ is } \mathcal{B}_X\text{-measurable}\}.$$

Then \mathcal{D} is a Dynkin class and, from Eq. (17), \mathcal{D} contains the class

$$\{B \times \Omega^\infty | B \in \mathcal{B}^n (1 \leq n < \infty)\},$$

which generates \mathcal{B}^∞ . Therefore, Dynkin's theorem assures that $\tilde{\nu}_x^\infty(\cdot)$ is a Markov kernel. Thus, we can define a POVM \tilde{E}_∞ by

$$\tilde{E}_\infty(B) := \int_{\Omega_X} \nu_x^\infty(B) F(dx) \quad (B \in \mathcal{B}^\infty),$$

which satisfies $\tilde{E}_\infty \leq F$. Then for each $B \in \mathcal{B}^n$, we have

$$\tilde{E}_\infty(B \times \Omega^\infty) = \int_{\Omega_X} \tilde{\nu}_x^n(B \times \Omega_X) F(dx) = (I^{*n} * F)(B \times \Omega_X) = E_\infty(B \times \Omega^\infty),$$

where we have used Eq. (13) in the second equality. This implies that $E_\infty = \tilde{E}_\infty \leq F$, which completes the proof. \square

The part of the result of Theorem 4 (conservation of E_∞ by I) was first obtained in the Ph.D. thesis by the author.¹⁷

IV. EXAMPLES OF THE INFINITE COMPOSITION: PHOTON COUNTING AND QUANTUM COUNTER MEASUREMENTS

In this section, we consider typical examples of standard Borel CP instruments, namely, photon counting^{1,3,18,19} and quantum counter^{3,8,20} instruments and evaluate the infinite compositions of them.

Let the system Hilbert space \mathcal{H} correspond to a single-mode photon field, and have a complete orthonormal system $\{|n\rangle\}_{n \in \mathbb{N}}$ called photon number eigenstates. Here the set of natural numbers \mathbb{N} contains 0. We denote the power set of \mathbb{N} as $2^{\mathbb{N}}$, and the countable product space of $(\mathbb{N}, 2^{\mathbb{N}})$ as $(\mathbb{N}^\infty, \mathcal{B}(\mathbb{N}^\infty))$.

The photon counting and quantum counter instruments for a finite time interval $t > 0$ are discrete and pure CP instruments with an outcome space $(\mathbb{N}, 2^{\mathbb{N}})$ defined by

$$\begin{aligned}I_B^{\text{pc}}(b) &:= \sum_{m \in B} M_m^{\text{pc}*} b M_m^{\text{pc}}, \\ M_m^{\text{pc}} &:= \sum_{n=0}^{\infty} \sqrt{p^{\text{pc}}(m+n|n)} |n\rangle \langle n+m|, \\ p^{\text{pc}}(m|n) &:= \binom{n}{m} (1 - e^{-\lambda t})^m e^{\lambda t(n-m)},\end{aligned}$$

for the photon counting instrument,^{1,3,18,21} and

$$\begin{aligned}\mathcal{I}_B^{\text{qc}}(b) &:= \sum_{m \in B} M_m^{\text{qc}*} b M_m^{\text{qc}}, \\ M_m^{\text{qc}} &:= \sum_{n=0}^{\infty} \sqrt{p^{\text{qc}}(m|n)} |n+m\rangle \langle n|, \\ p^{\text{qc}}(m|n) &:= \binom{n+m}{m} (e^{\lambda t} - 1)^m e^{-\lambda t(n+m+1)},\end{aligned}$$

for the quantum counter instrument.^{3,8} Here λ is a positive constant corresponding to the coupling strength between the detector and the photon field. The infinite composition of $\mathcal{I}^{\text{pc},\text{qc}}$ is a POVM $E_{\infty}^{\text{pc},\text{qc}}$ with the infinite product outcome space $(\mathbb{N}^{\infty}, \mathcal{B}(\mathbb{N}^{\infty}))$. Abusing the notation, $\mathcal{I}_{\{m\}}^{\text{pc},\text{qc}}(\cdot)$ is denoted as $\mathcal{I}_m^{\text{pc},\text{qc}}(\cdot)$.

We define the photon number observable $E^N(\cdot)$ by

$$E^N(B) := \sum_{n \in B} |n\rangle \langle n| \quad (B \in 2^{\mathbb{N}}),$$

and a POVM E^X with its outcome space $(\mathbb{R}_+, \mathcal{B}(\mathbb{R}_+))$, where \mathbb{R}_+ is a real half-line $(0, \infty)$ and $\mathcal{B}(\mathbb{R}_+)$ is the Borel σ -algebra of \mathbb{R}_+ , by

$$\begin{aligned}E^X(B) &= \int_B F_x dx \quad (B \in \mathcal{B}(\mathbb{R}_+)), \\ F_x &= \sum_{n \in \mathbb{N}} \frac{e^{-x} x^n}{n!} |n\rangle \langle n|.\end{aligned}\tag{18}$$

Here, dx in Eq. (18) is the ordinary Borel measure on the real line.

The following theorem gives explicit forms of the infinite compositions of $\mathcal{I}^{\text{pc},\text{qc}}$.

Theorem 5.

1. $E_{\infty}^{\text{pc}} \simeq E^N$,
2. $E_{\infty}^{\text{qc}} \simeq E^X$.

Proof.

1. From

$$\begin{aligned}\mathcal{I}_m^{\text{pc}}(|n_1\rangle \langle n_1|) &= p^{\text{pc}}(m|m+n_1) |m+n_1\rangle \langle m+n_1| \\ &= \sum_{n \in \mathbb{N}} \delta_{n,m+n_1} p^{\text{pc}}(m|n) |n\rangle \langle n|\end{aligned}\tag{19}$$

and

$$\begin{aligned}\sum_{m,n_1 \in \mathbb{N}} \delta_{n,m+n_1} \mathcal{I}_m^{\text{pc}}(|n_1\rangle \langle n_1|) &= \sum_{m,n_1 \in \mathbb{N}} \delta_{n,m+n_1} p^{\text{pc}}(m|m+n_1) |m+n_1\rangle \langle m+n_1| \\ &= |n\rangle \langle n|,\end{aligned}$$

we obtain $\mathcal{I}^{\text{pc}} * E^N \simeq E^N$. Thus, from Theorem 4, $E_{\infty}^{\text{pc}} \leq E^N$ holds.

In order to show $E^N \leq E_{\infty}^{\text{pc}}$, let us define $\mathcal{B}(\mathbb{N}^{\infty})$ -measurable stochastic variables $M_k(\mathbf{m}^{(\infty)})$ and $M_{\infty}(\mathbf{m}^{(\infty)})$ by

$$\begin{aligned}M_k(\mathbf{m}^{(\infty)}) &:= \sum_{i=1}^k m_i \in \mathbb{N}, \\ M_{\infty}(\mathbf{m}^{(\infty)}) &:= \lim_{k \rightarrow \infty} M_k(\mathbf{m}^{(\infty)}) \in \mathbb{N} \cup \{\infty\},\end{aligned}\tag{20}$$

where $\mathbf{m}^{(\infty)} := (m_1, m_2, \dots) \in \mathbb{N}^{\infty}$. Since $\delta_{m, M_k(\mathbf{m}^{(\infty)})} \rightarrow \delta_{m, M_{\infty}(\mathbf{m}^{(\infty)})}$ for each $m \in \mathbb{N}$ and each $\mathbf{m}^{(\infty)} \in \mathbb{N}^{\infty}$, we have

$$\begin{aligned} E^{M_\infty}(m) &:= \int_{\mathbb{N}^\infty} \delta_{m, M_\infty(m^{(\infty)})} E_\infty^{\text{pc}}(d\mathbf{m}^{(\infty)}) \\ &= \lim_{k \rightarrow \infty} \int_{\mathbb{N}^\infty} \delta_{m, M_k(m^{(\infty)})} E_\infty^{\text{pc}}(d\mathbf{m}^{(\infty)}) \\ &= \lim_{k \rightarrow \infty} \sum_{m_1, \dots, m_k \in \mathbb{N}} \delta_{m_1 + \dots + m_k, m} \mathcal{I}_{m_1}^{\text{pc}} \circ \dots \circ \mathcal{I}_{m_k}^{\text{pc}}(I). \end{aligned} \quad (21)$$

Here $E^{M_\infty}(m) = E^{M_\infty}(\{m\})$ is the POVM derived by M_∞ and the limit is in the sense of the weak operator topology. Let us evaluate Eq. (21). From Eq. (19), for each $k \geq 1$ and $(m_1, \dots, m_k, n_k) \in \mathbb{N}^{k+1}$ we have

$$\begin{aligned} &\mathcal{I}_{m_1}^{\text{pc}} \circ \dots \circ \mathcal{I}_{m_k}^{\text{pc}}(|n_k\rangle\langle n_k|) \\ &= \left(\prod_{i=1}^k p^{\text{pc}}(m_i | m_i + \dots + m_k + n_k) \right) |m_1 + \dots + m_k + n_k\rangle\langle m_1 + \dots + m_k + n_k|, \end{aligned}$$

and thus

$$\begin{aligned} &\mathcal{I}_{m_1}^{\text{pc}} \circ \dots \circ \mathcal{I}_{m_k}^{\text{pc}}(I) \\ &= \sum_{m_1, \dots, m_k, n \in \mathbb{N}} p_k^{\text{pc}}(m_1, \dots, m_k | n) |n\rangle\langle n|, \end{aligned}$$

where

$$p_k^{\text{pc}}(m_1, \dots, m_k | n) := \prod_{i=1}^k p^{\text{pc}}(m_i | n - m_1 - \dots - m_{i-1}). \quad (22)$$

By performing some calculations, the distribution of $m_1 + \dots + m_k$ for conditional distribution Eq. (22) is evaluated to be

$$\begin{aligned} p_k^{\text{pc}}(m | n) &:= \sum_{m_1, \dots, m_k \in \mathbb{N}} \delta_{m_1 + \dots + m_k, m} p_k^{\text{pc}}(m_1, \dots, m_k | n) \\ &= \binom{n}{m} (1 - e^{-\lambda t k})^m (e^{-\lambda t k})^{n-m} \\ &\rightarrow \delta_{n, m} \quad (k \rightarrow \infty). \end{aligned}$$

Thus, from Eq. (21), we have

$$E^{M_\infty}(m) = \lim_{k \rightarrow \infty} \sum_{n \in \mathbb{N}} p_k^{\text{pc}}(m | n) |n\rangle\langle n| = |m\rangle\langle m|, \quad (23)$$

which implies that $E^N = E^{M_\infty} \leq E_\infty^{\text{pc}}$ and we have proved the assertion. Note that Eq. (23) indicates that $E^{M_\infty}(\infty) = I - E^{M_\infty}(\mathbb{N}) = O$, i.e., M_k is convergent E_∞^{pc} -almost surely.

2. From

$$\begin{aligned} \mathcal{I}_m^{\text{qc}}(F_x) &= e^{-\lambda t} p^{\text{qc}}(m | e^{-\lambda t} x) F_{e^{-\lambda t} x}, \\ p^{\text{qc}}(m | x) &:= \frac{[(e^{\lambda t} - 1)x]^m}{m!} \exp[-(e^{\lambda t} - 1)x], \end{aligned}$$

we have

$$\begin{aligned} \mathcal{I}_{m_1}^{\text{qc}} \circ \dots \circ \mathcal{I}_{m_k}^{\text{qc}}(F_x) &= e^{-\lambda t k} \left(\prod_{i=1}^k p^{\text{qc}}(m_i | e^{-\lambda t(k-i+1)} x) \right) F_{e^{-\lambda t k} x}, \\ \mathcal{I}_{m_1}^{\text{qc}} \circ \dots \circ \mathcal{I}_{m_k}^{\text{qc}}(I) &= \int_0^\infty \mathcal{I}_{m_1}^{\text{qc}} \circ \dots \circ \mathcal{I}_{m_k}^{\text{qc}}(F_x) dx \\ &= \int_0^\infty \left(\prod_{i=1}^k p^{\text{qc}}(m_i | e^{\lambda t(i-1)} x) \right) E^X(dx). \end{aligned} \quad (24)$$

Let $\nu_x^\infty(\cdot)$ be the product measure of $p^{\text{qc}}(\cdot|e^{\lambda t(i-1)}x)$ with respect to $i \geq 1$ with its outcome space $(\mathbb{N}^\infty, \mathcal{B}(\mathbb{N}^\infty))$. Then Eq. (24) implies that

$$E_\infty^{\text{qc}}(\cdot) = \int_0^\infty \nu_x^\infty(\cdot) E^X(dx), \quad (25)$$

and we have shown $E_\infty^{\text{qc}} \leq E^X$. To show $E^X \leq E_\infty^{\text{qc}}$, let us define $\mathcal{B}(\mathbb{N}^\infty)$ -measurable stochastic non-negative variables $X_k(\mathbf{m}^{(\infty)})$ by

$$X_k(\mathbf{m}^{(\infty)}) := e^{-\lambda tk} \sum_{i=1}^k m_i. \quad (26)$$

If we denote the expectation with respect to $\nu_x^\infty(\cdot)$ as $\mathbb{E}_x[\cdot]$, we have

$$\begin{aligned} \mathbb{E}_x[X_k] &= (1 - e^{-\lambda tk})x, \\ \mathbb{E}_x[(X_k - \mathbb{E}_x[X_k])^2] &= e^{-\lambda tk}(1 - e^{-\lambda tk})x, \end{aligned}$$

where we have omitted the dependence of $\mathbf{m}^{(\infty)} \in \mathbb{N}^\infty$. Thus, from Chebyshev's inequality we obtain

$$\nu_x(\{|X_k - x| > e^{-\lambda tk/4}\}) \leq e^{\lambda tk/2} \mathbb{E}_x[|X_k - x|^2] \leq C_x e^{-\lambda tk/2},$$

where C_x is some positive constant independent of k . Then Borel-Cantelli lemma assures that X_k converges to x ν_x^∞ -almost surely. Therefore, from Eq. (25), X_k is convergent E_∞^{qc} -almost surely. Then, we can define a non-negative stochastic variable $X_\infty := \lim_{k \rightarrow \infty} X_k$, which satisfies $X_\infty = x$ ν_x^∞ -almost surely, i.e., $\nu_x(X_\infty^{-1}(B)) = \chi_B(x)$ for each $B \in \mathcal{B}(\mathbb{R}_+)$. Therefore, for each $B \in \mathcal{B}(\mathbb{R}_+)$, we have

$$\begin{aligned} \int_{\mathbb{N}^\infty} \chi_B(X_\infty(\mathbf{m}^{(\infty)})) E_\infty^{\text{qc}}(d\mathbf{m}^{(\infty)}) &= E_\infty^{\text{qc}}(X_\infty^{-1}(B)) \\ &= \int_0^\infty \nu_x^\infty(X_\infty^{-1}(B)) E^X(dx) \\ &= \int_0^\infty \chi_B(x) E^X(dx) = E^X(B), \end{aligned}$$

which implies $E^X \leq E_\infty^{\text{qc}}$. Thus, we have proved $E_\infty^{\text{qc}} \simeq E^X$. \square

According to the above proof, we also obtain the following theorem concerning the convergences of the stochastic variables M_k and X_k .

Theorem 6.

- (i) M_k defined by Eq. (20) is convergent E_∞^{pc} -almost surely and the POVM corresponding to the distribution of $\lim_{k \rightarrow \infty} M_k$ coincides with the photon number observable E^N .
- (ii) X_k defined by Eq. (26) is convergent E_∞^{qc} -almost surely and the POVM corresponding to the distribution of $\lim_{k \rightarrow \infty} X_k$ coincides with E^X .

We remark that the statement of Theorem 6 (ii) is essentially the same as Theorem 4 of Ref. 8 while the proof of Ref. 8 is, rigorously speaking, insufficient due to the following reason. The authors of Ref. 8 show that the characteristic function of X_k converges to that of E^X and conclude the assertion of the theorem. However, as well-known in the measure theoretic probability theory, the pointwise convergence of the characteristic function, which is equivalent to the convergence in distribution, does not necessarily imply the almost sure convergence of a stochastic variable. In this sense, our proof of Theorem 6 (ii) complements the mathematically insufficient discussion of Ref. 8.

V. SUMMARY

In summary, we have rigorously reformulated the concept of the information conservation condition in Definition 1 depending on the equivalence relation among POVMs and the composition

between an instrument and a POVM. By using quantum Kolmogorov extension theorem, we have constructed the infinite composition of a given standard Borel CP instrument \mathcal{I} . We have shown that the infinite composition is the least informative standard Borel POVM that is conserved by \mathcal{I} . We have considered specific examples of CP instruments, namely, photon counting and quantum counter instruments, and shown that their infinite compositions are equivalent to the photon number observable E^N and the POVM E^X given by Eq. (18), respectively. As a by-product of the proof, we have found some results on the almost sure convergences of the total counting number for the photon counting and the properly normalized counting number for the quantum counter cases, respectively. The latter result on the convergence in the quantum counter measurement complements the insufficiency of the proof in the existing work⁸ from the standpoint of the rigorous measure theoretic description of quantum measurements.

ACKNOWLEDGMENTS

The author acknowledges supports by Japan Society for the Promotion of Science (KAKENHI Grant No. 269905). He also would like to thank helpful discussions with Masahito Ueda (the University of Tokyo) and Tomohiro Shitara (the University of Tokyo).

- ¹ M. Srinivas and E. Davies, "Photon counting probabilities in quantum optics," *Opt. Acta* **28**, 981–996 (1981).
- ² M. Ozawa, "Quantum measuring processes of continuous observables," *J. Math. Phys.* **25**, 79–87 (1984).
- ³ Y. Kuramochi and M. Ueda, "Classicality condition on a system observable in a quantum measurement and a relative-entropy conservation law," *Phys. Rev. A* **91**, 032110 (2015).
- ⁴ S. Dorofeev and J. de Graaf, "Some maximality results for effect-valued measures," *Indag. Math. (N.S.)* **8**, 349–369 (1997).
- ⁵ W. M. de Muynck, *Foundations of Quantum Mechanics, an Empiricist Approach* (Kluwer Academic Publishers, Dordrecht, Boston, London, 2002).
- ⁶ T. Heinonen, "Optimal measurements in quantum mechanics," *Phys. Lett. A* **346**, 77–86 (2005).
- ⁷ A. Jenčová, S. Pulmannová, and E. Vinceková, "Sharp and fuzzy observables on effect algebras," *Int. J. Theor. Phys.* **47**, 125–148 (2008).
- ⁸ M. Ueda, N. Imoto, and H. Nagaoka, "Logical reversibility in quantum measurement: General theory and specific examples," *Phys. Rev. A* **53**, 3808–3817 (1996).
- ⁹ M. Ban, "State reduction, information and entropy in quantum measurement processes," *J. Phys. A: Math. Gen.* **32**, 1643 (1999).
- ¹⁰ P. R. Halmos and L. J. Savage, "Application of the Radon-Nikodym theorem to the theory of sufficient statistics," *Ann. Math. Stat.* **20**, 225–241 (1949).
- ¹¹ T. Heinosaari and M. Ziman, *The Mathematical Language of Quantum Theory: From Uncertainty to Entanglement* (Cambridge University Press, 2011).
- ¹² E. B. Davies, *Quantum Theory of Open Systems* (IMA, 1976).
- ¹³ E. B. Davies and J. Lewis, "An operational approach to quantum probability," *Commun. Math. Phys.* **17**, 239–260 (1970).
- ¹⁴ S. M. Srivastava, *A Course on Borel Sets* (Springer, 1998).
- ¹⁵ R. Tumulka, "A Kolmogorov extension theorem for POVMs," *Lett. Math. Phys.* **84**, 41–46 (2008).
- ¹⁶ A. N. Kolmogorov, *Grundbegriffe der Wahrscheinlichkeitsrechnung* (Springer, 1933).
- ¹⁷ Y. Kuramochi, "Relative-entropy conservation law in quantum measurement and its applications to continuous measurements," Ph.D. thesis, Department of Physics, The University of Tokyo, 2015, to be published in 2016 in http://repository.dl.itc.u-tokyo.ac.jp/index_e.html.
- ¹⁸ M. Ueda, "Probability-density-functional description of quantum photodetection processes," *Quantum Opt.: J. Eur. Opt. Soc. Part B* **1**, 131 (1989).
- ¹⁹ N. Imoto, M. Ueda, and T. Ogawa, "Microscopic theory of the continuous measurement of photon number," *Phys. Rev. A* **41**, 4127–4130 (1990).
- ²⁰ M. Ueda and M. Kitagawa, "Reversibility in quantum measurement processes," *Phys. Rev. Lett.* **68**, 3424–3427 (1992).
- ²¹ M. Ueda, N. Imoto, and T. Ogawa, "Quantum theory for continuous photodetection processes," *Phys. Rev. A* **41**, 3891–3904 (1990).